



A new relaxed PSS preconditioner for nonsymmetric saddle point problems[☆]



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ABSTRACT

A new relaxed PSS-like iteration scheme for the nonsymmetric saddle point problem is proposed. As a stationary iterative method, the new variant is proved to converge unconditionally. When used for preconditioning, the preconditioner differs from the coefficient matrix only in the upper-right components. The theoretical analysis shows that the preconditioned matrix has a well-clustered eigenvalues around (1, 0) with a reasonable choice of the relaxation parameter. This sound property is desirable in that the related Krylov subspace method can converge much faster, which is validated by numerical examples.

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1. Introduction

The steady-state Navier–Stokes system is a basic tool for the modeling of an incompressible Newtonian fluid [30]. Finding effective methods for this system remains a crucial issue for a range of engineering applications [25,26]. Let $\Omega \subset \mathbb{R}^2$ (or \mathbb{R}^3) be a bounded, connected domain with a boundary Γ . Given a force field \mathbf{f} and boundary data \mathbf{g} , the problem is to ascertain a velocity field \mathbf{u} and a pressure field \mathbf{p} that satisfy the steady-state Navier–Stokes system [15]

$$\begin{aligned} -\nu \Delta \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \nabla \mathbf{p} &= \mathbf{f}, & \text{in } \Omega, \\ \operatorname{div} \mathbf{u} &= 0, & \text{in } \Omega, \\ \mathbf{u} &= \mathbf{g}, & \text{on } \partial\Omega, \end{aligned} \quad (1)$$

where $\nu > 0$ is the kinematic viscosity, Δ is the vector Laplacian, ∇ is the gradient and div is the divergence. The Navier–Stokes Eq. (1) is nonlinear due to the existence of $(\mathbf{u} \cdot \nabla) \mathbf{u}$. Two common strategies, namely Picard iteration and Newton iteration, are available for linearizing (1), which leads to the following Oseen problem

$$\begin{aligned} -\nu \Delta \mathbf{u} + (\boldsymbol{\omega} \cdot \nabla) \mathbf{u} + \nabla \mathbf{p} &= \mathbf{f}, & \text{in } \Omega, \\ \operatorname{div} \mathbf{u} &= 0, & \text{in } \Omega, \\ \mathbf{u} &= \mathbf{g}, & \text{on } \partial\Omega, \end{aligned} \quad (2)$$

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